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SOLAR EXTINCTION RADIOMETRY

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- I. During this period the work on the spectral line parameters of the OR $1^2\Sigma = X^2\Pi(0.0)$ band has been completed. A detailed paper describing this work has been published in J.Q.S.R.T. and a reprint is enclosed as an Appendix. Numerous reprints requests for this paper were received.
- II. Using the methods developed for NO_2 quantification (1) during previous phases of this project, the uv-visible data obtained during the 1977 balloon flights have been used for O_3 quantification.

For etmospheric ozone amounts and the region covered by these spectra, the ozone absorption renders the atmosphere opaque at low sun agles in the 2809-3100% region (Hartley bands). From about 3500 to 4000% the ozone absorption coefficients are too small to give appreciable absorption even at the lowest sun angles scanned (-92). From 4000 to 5000% (Chappuis bands) the atmospheric absorption is quite weak and broad and does not have the contrast structure required for the three wavelengths method. The region from 3100% to 3500% appears to be the best region to use for determining ozone columns with the three-wavelength method.

Accurate determination of ozone columns and volume mixing ratios requires accurate ozone absorption coefficients measured at atmospheric temperatures (210 to 260K). The best presently available published absorption coefficients are those of Vigroux². We have used these coefficients and five combinations of three wavelengths chosen from the features marked in Figure 1

to determine azone columns and volume mixing ratios.

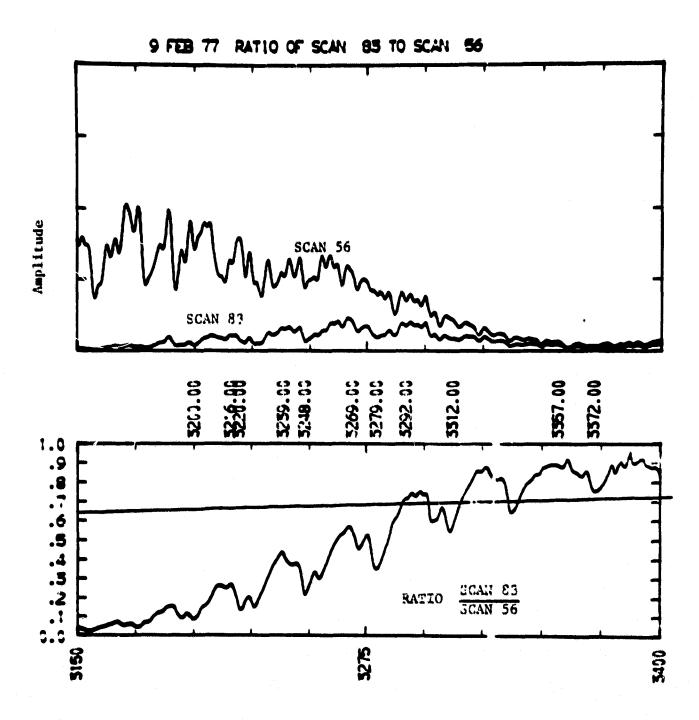
The ozone volume mixing ratios determined from the 9 February 1977 UV and 17 February 1977 UV data have been compared with standard middle-latitude ozone profiles³. The ozone profiles from both flights are in agreement with the standard profile mixing ratios, as shown in Figure 2.

The spectra in Fig. 1 have been smoothed to 2.51 resolution. It should be noted that while the fine structure of the atmospheric Ω_2 absorption in Fig. 1 is not fully accounted for the Vigroux coefficients, it is in agreement with the recent high resolution (0.2%) work being completed at the N.P.S. The accuracy of the Ω_2 profiles should improve when this work, which is conducted at atmospheric temperatures, is completed.

III. Work has started on the new UV solar spectra obtained during the 3/19/81 balloon flight. Numerous high and low sun scans were obtained during ascent and from float altitude (~33 km) at 0.03% resolution in the 3063-3099% region. High and low sun scans from float altitude have been calibrated in wavelength. Representative scans are shown in Fig. 3 and Fig. 4 respectively for the 3075-3082% region, which is considered best for OH measurements. The spectra are being studied for OH identification and quantification.

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WAVELENGTH (Å)

Fig. 1 Ratio of Scan 83 (low sun) to Scan 56 (high sun). The fine structure in the ratio spectrum is due to atmospheric 0_3 . Resolution is 2.5\AA .

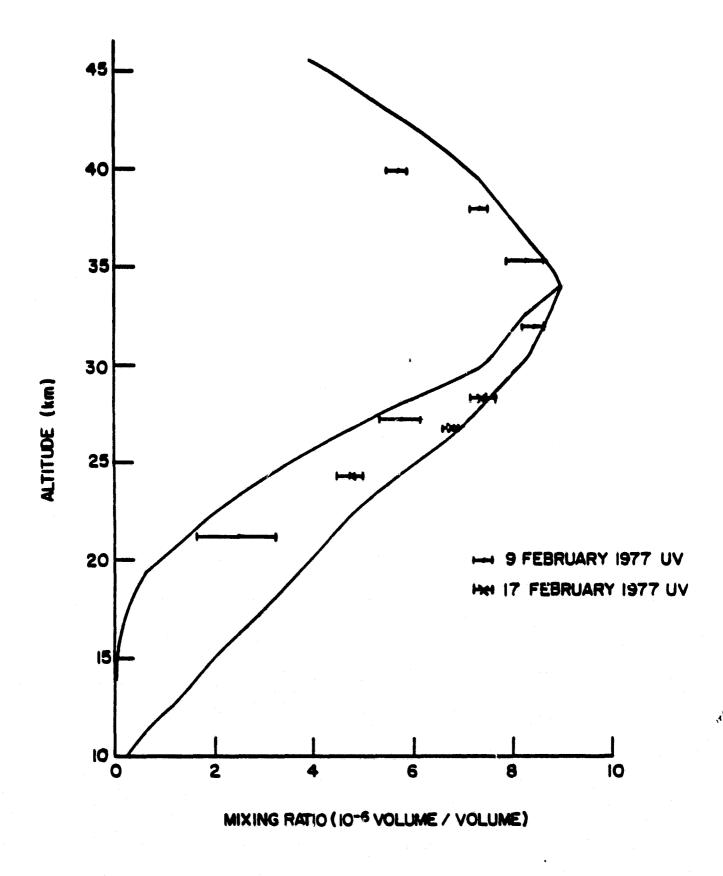


Figure 2 Preliminary czone profiles superimposed on the standard midlatitude ozone profile of Mateer et al. (1980).

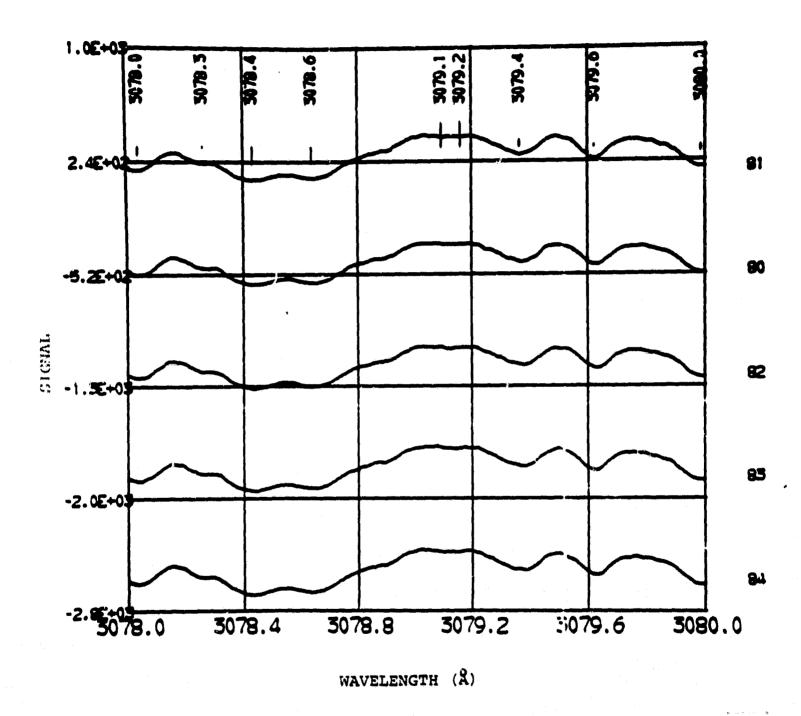


Fig. 3a Selected high sun scans obtained during the 3/19/81 balloon flight from float altitude (33 km). Resolution is 0.03Å.

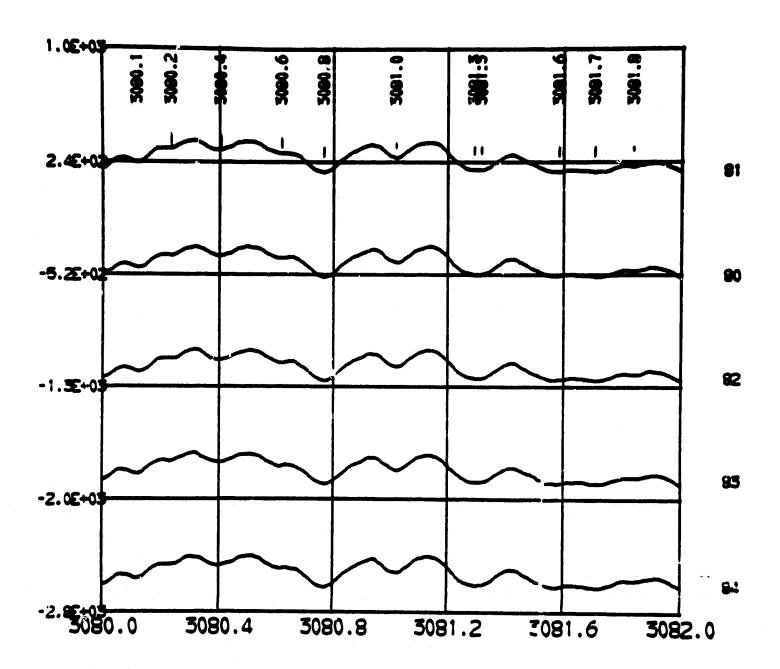


Fig. 3b

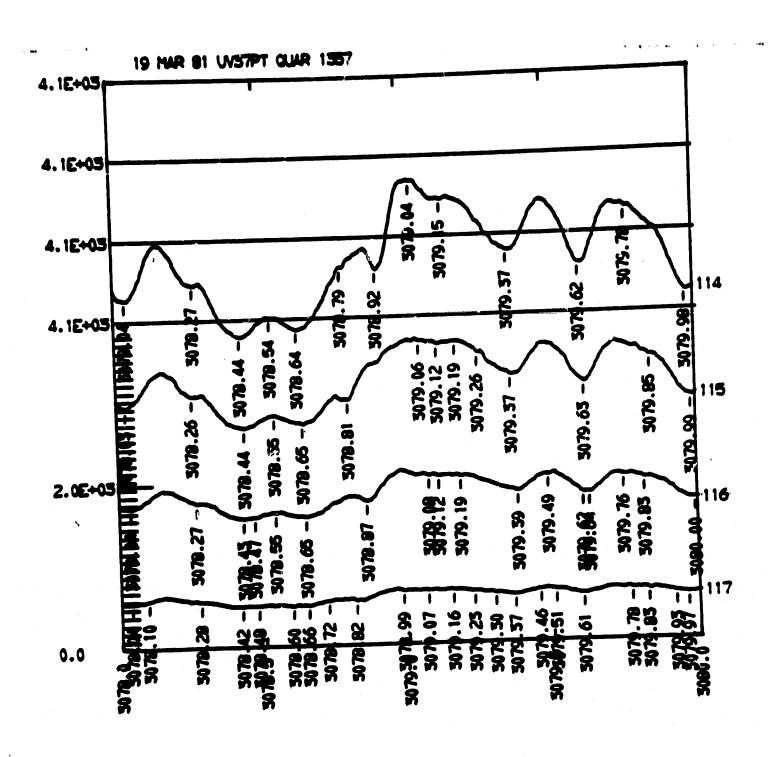


Fig. 4a Selected low sun scans obtained during the 3/19/81 balloon flight from float altitude (33 km). Resolution is 0.03Å.

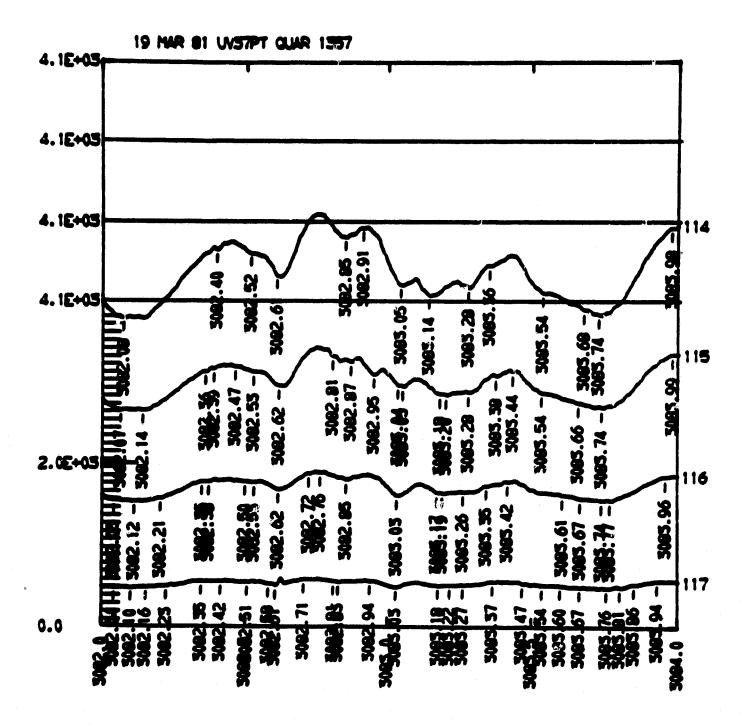


Fig. 4b

SPECTRAL LINE PARAMETERS FOR THE A'\(\Sigma - X'\) II(0, 0) BAND OF OH FOR ATMOSPHERIC AND HIGH TEMPERATURES

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Abstract—Individual spectral line parameters including line positions, strengths, and intensities, have be: if generated for the $A^2\Sigma-X^2\Pi(0,0)$ band of OH, applicable to atmospheric and high temperatures. Energy levels and transition frequencies are calculated by numerically disponalizing the Hamiltonian. Line strengths are calculated using the displayment and organization derived from energy matrix diagonalization. The line strengths are compared to those calculated from previously published algebraic line strength formulas. Tables of line parameters are presented for 240 and 4400°K.

I. INTRODUCTION

The $A^2\Sigma - X^2\Pi(0,0)$ band of OH in the 3085 Å region has been of interest to quantitative spectroscipists for many years because of its high absorption and emission intensity and convenient wavelength location for spectroscopic probes. The hydroxyl radical is a common by-product of most combustion processes, is present that it is atmospheric, solar and stellar spectra, and in recent years has been also recognized as an important trace constituent in atmospheric chemistry. Accurate determination of the amount of OH present during spectroscopic experiments depends on precise knowledge of line positions and intensities. Several analyses of spectral line positions for this band are available; among the more important of these are those of Dieke and Crosswhite, who provided the first extensive analysis of the OH u.v. spectrum, and Destombes et al., who performed elaborate analysis of modern microwave, i.r. and u.v. OH data. Intensity (relative and absolute) studies of this band have been reviewed recently by Chidsey and Crosley, who also performed extensive RKR calculations of rotational transition probabilities for the A-X system of OH.

The purpose of this work is to combine the best presently available data and theory to derive accurate quantitative line parameters for the $A^2\Sigma - X^2\Pi(0,0)$ band, applicable to atmospheric and high temperatures. The results are displayed in line parameter tables, and include improved values for the line strength, calculated in intermediate coupling from the energy matrix eigenvectors.

2. LINE PARAMETERS DERIVATION

The OH molecule has an unpaired electron with total electronic angular momentum L=1 and S=1/2. In the electronic ground state the projection of L along the internuclear axis is $\Lambda=\pm 1$. The projection of S along the internuclear axis is $\Sigma=\pm 1/2$, with a total electronic angular momentum projection $\Omega=\Lambda+\Sigma$. Here, Λ , Σ , Ω are considered as signed quantities, as in the notation of Hougen.⁴ The electronic ground state is an inverted ${}^2\Pi$ state with the ${}^2\Pi_{1/2}(\Omega=\pm 1/2,\,F_2)$ levels at higher energy than the ${}^2\Pi_{3/2}$ levels $(\Omega=\pm 3/2,\,F_1)$. The rotational levels for this state are intermediate between Hund's cases (a) and (b). In the ${}^2\Sigma$ upper state, which is Hund's case (b), $\Lambda=0$ and $\Omega=\pm 1/2$ with $J=N\pm 1/2$. The ${}^2\Pi$ and ${}^2\Sigma$ states perturb one another and produce Λ doubling for each N (${}^2\Sigma_{1/2}$) or J (${}^2\Pi_{1/2,3/2}$).

We use the unique perturber approximation described by Destombes et al.² to calculate energy levels. This process is restricted to a single vibrational level v in the $A^2\Sigma - X^2\Pi$ subspace. The total angular momentum number F (not to be confused with the level designations F_1 and F_2) is a good quantum number. For a given F, the F and F levels are weakly coupled by magnetic hyperfine interactions. These interactions are negligible in the calculation of electronic spectra and F may be considered a good quantum number. This procedure gives a F hamiltonian matrix, the elements of which are listed in Table 1. The matrix elements are

Table 1. Nonzero Hamiltonian matrix elements (Hund's case (a) basis).

Metrix Clearat	Yelve
«283/2 [N] 283/2"	$B_{\frac{1}{2}}(\pi^{2}-2)+\frac{h}{2}-ij_{\frac{1}{2}}(\pi^{4}-3\pi^{2}+3)+h_{\frac{1}{2}}(\pi^{2}-2)+ij_{\frac{1}{2}}(\pi^{6}-3\pi^{4}+5\pi^{2}-6)$
48 1 N 381/2	$B_{R} n^{2} - \frac{h}{2} = B_{R} (n^{4} + n^{2} - 1) - A_{D} n^{2} + M_{R} (n^{6} + 3n^{4} - 5n^{2} + 2)$
4281/2 N 283/22	B _N y-2D _Q y(x ² -1)+M _Q y(3x ⁶ -5x ² +3)
4251/2 0 251/25	$3e^{\frac{1}{2}\sqrt{\frac{1}{6}}}e^{+\sqrt{3}} + 0e^{(\pi^{3}+\pi^{2})+6}e^{(\pi^{6}+3\pi^{6})}$
4251/2 N 25-1/2>	B _E x-2B _E x ³ 4H _E (3x ⁵ 4x ³)
$e^2\pi_{1/2} u ^{-2}\pi_{1/2}$	$(8L_{\phi})y = (0L_{\phi})y(2\pi^{2}-1) + (\Delta_{\phi}L_{\phi})y + (8L_{\phi})y(3\pi^{4}+1)$
$e^2\pi_{1/2} \mathbf{u} ^{-2}\pi_{-1/2}$	${\rm est}_{a}{\succ}{\rm m-eod}_{a}{\succ}{\rm m}(2{\rm m}^2{+}1) \leftarrow {\rm edd}_{a}{\succ}{\rm m}(3{\rm m}^4{\leftarrow}6{\rm m}^2{-}2)$
<283/2 M 25,1/2>	-2 <dl<sub>p>my+3<ml<sub>p>m(2m²-1)y</ml<sub></dl<sub>
4251/2 M 251/2"	<8L_>++++++++++++++++++++++++++++++++++++

Xetes:

- 1. =J+1/2; $y=((J-1/2)(J+3/2))^{1/2}$
- Hatrim elements are unchanged by anchange of initial and final states or by setting 0 to -0 in both initial and final states.

written in Hund's case (a), with the wave functions represented by $|\Delta S\Sigma\rangle|J\Omega\rangle = |\Delta S\Sigma;J\Omega\rangle$, so that

$${}^{2}\Sigma_{\pi 1/2}^{2}: |\Delta S\Sigma; J\Omega\rangle = |0^{+}1/2 \pm 1/2; J \pm 1/2\rangle,$$

$${}^{2}\Pi_{\pi 1/2}^{2}: |\Delta S\Sigma; J\Omega\rangle = |\pm 1|1/2 \pm 1/2; J \pm 1/2\rangle,$$

$${}^{2}\Pi_{\pi 3/2}^{2}: |\Delta S\Sigma; J\Omega\rangle = |\pm 1|1/2 \pm 1/2; J \pm 3/2\rangle.$$
(1)

The Hamiltonian constants used are shown in Table 2 (these constants give a better fit to the

Table 2. Hamiltonian constants for A2Z-X2II(0,0).

Constant	Value * (cm=1)
10	16.9258478
D ₂	2.0396602 x 10"
N.	97.7612 × 10"
Yg	-7.8399 × 10 ⁻³
v _a	32402.056239
A.	-139.228125
3.	18.5497354
0,	1.907es2 x 30 ^{m3}
W ₃	0.1239a16 # 10"6
A 5	-0.7233. x 10 ⁻³
«AL ₊ »	-151.9226:12
<8L_>	25.0435440
<dl_></dl_>	2,6923336 × 10 ⁻¹
«A ₃ L ₄ »	8.651427 x 10 ⁻³
4E.>	0.166gg x 10-6

^{*} Values are from Faf. 5 and are rounded to three figures beyond the standard errors indicated in Ref. 2.

observed spectrum than those of Ref. 2, which do not have a sufficient number of digits retained). The 6×6 Hamiltonian matrix may be reduced to two 3×3 blocks by the Kronig transformation.

$$|J\Omega\delta\rangle = \frac{1}{2}(|\Lambda S\Sigma\rangle|J\Omega\rangle + \delta|-\Lambda S - \Sigma\rangle|J-\Omega\rangle), \tag{2}$$

where 8 equals s = symmetric = + or a = antisymmetric = -. This notation ε hould not be confused with the s, a and +, - notation for homonuclear diatomic molecules.† In this new basis, one of the 3×3 blocks contains only matrix elements of the type $\langle J'\Omega's|H|J\Omega s\rangle$ and the other contains only matrix elements of the type $\langle J'\Omega'a|H|J\Omega a\rangle$.

After numerical diagonalization of a 3×3 Kronig transformed block, the electronic state of each eigenvalue (energy level or, more properly, term value in cm⁻¹) may be determined by noting that the largest eigenvalue belongs to the $^2\Sigma_{1/2}$ state, the intermediate eigenvalue to the $^2\Pi_{1/2}$ state, and the smallest eigenvalue to the $^2\Pi_{3/2}$ state. The remaining quantum numbers and partities may be determined from Table 3.

The selection rules for the $A^2\Sigma - X^2\Pi(0,0)$ electric dipole allowed transitions are $\Delta J = 0, \pm 1, \Delta N = 0, \pm 1, \pm 2$. Parity selection rules require $+ \leftrightarrow -, + \leftrightarrow +,$ and $- \leftrightarrow -,$ which give for the wavefunction the Kronig symmetry selection rules of $s \leftrightarrow s$ and $a \leftrightarrow a$ for $\Delta J = \pm 1$ (P and R branches) and $s \leftrightarrow a$ for $\Delta J = 0$ (Q branches). Transitions are designated by $^{\Delta N}\Delta J_{PP}(J^P)$ where refers to the upper state and "refers to the lower state and F is F_1 or F_2 . These selection rules permit 12 branches, of which 6 are main branches $(^PP_{11}, ^QQ_{11}, ^RR_{11}, ^PP_{22}, ^QQ_{22}, ^RR_{22})$ and 6 are (weaker) satellite branches $^QP_{21}, ^RQ_{21}, ^RR_{21}, ^PQ_{12}, ^QR_{12}, ^PP_{12})$.

The line intensities $S_{le}(T)$ (cm⁻¹/atm-cm) at temperature $T(^{\circ}K)$ are calculated from

$$S_{lo}(T) = \frac{1}{8\pi c \nu^2} \left(\frac{N}{\rho} \right) \frac{e^{-1.4366g^2/T}}{Q_{ob}} A_i^{(2)} (2J' + 1)(1 - e^{-1.4366\pi^2}), \tag{3}$$

where $\nu (= E' - E'')$ is the transition frequency in cm⁻¹, $c = 2.99792458 \times 10^{10}$ cm/sec, N is the total number of OH molecules/cm³, p is the pressure in atm, $A_i^{\mu}J_i$ is the Einstein A coefficient in sec⁻¹, E'' is the lower state energy $(= E_i)$, and Q_{eR} is the vibration rotation partition function.

Line intensities may be converted from cm^{-1}/atm -cm at T to cm/molecule at T by

$$S_{lu}(T)(\text{cm/molecule}) = 3.721963 \times 10^{-20} \frac{T(^{\circ}\text{K})}{273.16(^{\circ}\text{K})} S_{lu}(T)(\text{cm}^{-1}/\text{cm-atm}).$$
 (4)

The $S_{le}(T)$ in cm/molecule are at the population temperature.

We assume that $Q_{\nu R}$ is given by $Q_{\nu}Q_{R}$, where the vibrational partition function Q_{ν} in the harmonic oscillator approximation is

$$Q_r = \frac{1}{1 - e^{-1.03\pi\omega_r/\gamma}}$$
 (5)

Table 3. Assignment of quantum numbers to eigenvalues for given J.

State	* R9/2	* A 1/2	122	/2
7	1	2	1	2
N	J- ¹ /2	3+1/2	$j_{-}^{1}/_{2}$	3+1/2
Parity: Eigonvalue from	•			
Symmetric block	(-1)***1	(-1) ^N	_	(-1) ^N
Anthoyumstric block	(-1) ^N	(-1) ^{N+1}	(-1) ^M	

^{*}We thank J. T. Hougen for pointing out that our notation s = symmetric = + is equivalent to the f levels and a = antisymmetric = - is equivalent to the e levels of J. M. Brown, J. T. Hougen, K.-P. Huber, J. W. C. Johns, I. Kopp, H. Lefebvre-Brion, A. J. Morer, D. A. Ramsay, J. Rostas, and R. N. Zare, J. Molec. Spectrosc, 56, 500 (1975).

and ω_r is the vibrational harmonic oscillator frequency in cm⁻¹. Huber and Herzberg⁷ give $\omega_r = 3737.76_1 \text{ cm}^{-1}$. A comparison of the values of Q_r calculated by Eq. (5) and by direct summation of $e^{-E_{cr}T}$ shows a difference of less than 0.2% at 4600°K. The rotational partition function Q_R is calculated from the actual energy levels as follows:

$$Q_R = \sum_{I \in I} (2J'' + 1) e^{-1.4300E^*(J^*\Pi^*)/T}.$$
 (6)

Chidsey and Crosley give the Einstein A coefficient as

$$A_{ij}^{\nu,j} = \frac{64\pi^4}{3h} \rho_{ij}^{\nu,j} S_{ij} \cdot \nu^3 / (2J' + 1) \sec^{-1}, \tag{7}$$

where $p \in \mathcal{T}$ is the rovibrational transition probability and S_{PP} is the rotational line strength. Chidsey and Crosley tabulate relative values of A_{PP}^{**} through N'' = 32 for the $A^2\Sigma - X^2\Pi(0,0)$ band in their Table 4. They state that they calculated the line strengths S_{PP} based on Earls' formulas with a J-dependent spin-orbit coupling parameter A and the rotational constants of Dieke and Crosswhite. We have found that the use of Earls' formulas can lead to significant errors at high J in the satellite bands. (This problem is further discussed in the following section.) Therefore, we have calculated S_{PP} by following the method described by Hougen. Chidsey and Crosley have kindly provided a table of relative values for $p \in \mathcal{T}$ through J = 35.5 prior to publication.

The A_{ij}^{eff} may be put on an absolute basis by noting that the lifetime of a state is

$$\tau_{v'J'} = \left(\sum_{v'J'} A_{v'J'}^{v'J'}\right)^{-1} \operatorname{sec},\tag{6}$$

Because Chidsey and Crosley³ give $A_1^0/A_0^0 \approx 0.0040$ (here the notation is A_v^{μ}), we assume that for the v'=0 vibrational state all vibrational states other than v''=0 make negligible contributions to $\tau_{v'J'}$. The best available lifetime for the rotationless (N'=0)v'=0 state is probably that measured by German, $\tau_{0,1/2}=(0.688\pm0.007)\times10^{-6}$ sec. There are three transitions from the v''=0 state to the rotationless v'=0 state, namely, $\tau_{0,1/2}=(0.5)$, and $\tau_{0,1/2}=(0.5)$. We calculate relative $\tau_{0,1/2}=(0.5)$ for all $\tau_{0,1/2}=(0.5)$ of interest by ignoring all constant factors in Eq. (7) and normalize these by using Eq. (8).

We form the S_{PP} by following the method given by Hougen.⁴ In addition to the selection rules $\Delta J=0, \pm 1$ and $+\leftrightarrow-$ parity, electric dipole selection rules on Ω give nonzero matrix elements only for $\langle \Omega\pm 1|\mu_z\pm i\mu_r|\Omega\rangle$ for $\Delta\Omega=\pm 1$ and $\langle \Omega|\mu_z|\Omega\rangle$ for $\Delta\Omega=0$, where μ_z , μ_r , and μ_r are electric dipole moment components in the molecule fixed axis system. In the laborate \sim axed coordinate system

$$\mu_{Z} = \frac{1}{2}(\alpha_{Zx} - i\alpha_{Zy})(\mu_{x} + i\mu_{y}) + \frac{1}{2}(\alpha_{Zx} + i\alpha_{Zy})(\mu_{x} - i\mu_{y}) + \alpha_{Zy}\mu_{z}, \tag{9}$$

where α_{Z_1} , α_{Z_2} , and α_{Z_2} are the direction cosines between the molecule-fixed and laboratory-

Table 4. Nonzero rotational direction cosine matrix elements in $(J'\Omega'|\mu_Z|J,\Omega)$ [Hund's case (a)].

	<u>J'=J+Ł</u>	<u>j'=j</u>	J'=J-1
<1;n u _z 1,n>	$\left[\frac{(J+\Omega+1)(J-\Omega+1)}{2(J+1)}\right]^{1/2}$	$-a \left[\frac{2J+1}{2J(J+1)} \right]^{1/2}$	((J+0) (J-0)) 1/2
<1;0+1 m 11:0>	$-\left[\frac{(J+0+1)(J+0+2)}{2(J+1)}\right]^{1/2}$	$\left[\frac{(J-\Omega)(J+\Omega+1)(2J+1)}{2J(J+1)}\right]^{1/2}$	$\left[\frac{(J-\Omega)(J-\Omega-1)}{2J}\right]^{1/2}$
<j;n-1 u<sub>2 /.n></j;n-1 u<sub>	$\left(\frac{(J-\Omega+1)(J-\Omega+2)}{2(J+1)}\right)^{1/2}$	$\left[\frac{(J+0)(J-0+1)(2J+1)}{2J(J+1)}\right]^{1/2}$	-((J+0) (J+0-1)) 1/2

fixed coordinates systems. The direction cosine matrix elements are given in Table 4. The μ_z , μ_y , and μ_z are taken to be experimentally determined parameters. For lack of better information, we take $\sqrt{(1/2)|\mu_x + i\mu_y|} = \sqrt{(1/2)|\mu_x - i\mu_y|} = \mu_x = 1$.

Because we assume that the electronic and rotational parts of the wavefunction are separable, the electric dipole matrix elements may be written as

$$\langle \Lambda'S'\Sigma'; J'\Omega'|\mu_{Z}|\Lambda S\Sigma; J\Omega \rangle = \frac{1}{2} \langle \Lambda'S'\Sigma'|\mu_{x} + i\mu_{y}|\Lambda S\Sigma \rangle \langle J'\Omega'|\alpha_{Zx} - i\alpha_{Zy}|J\Omega \rangle + \frac{1}{2} \langle \Lambda'S'\Sigma'|\mu_{x} - i\mu_{y}|\Lambda S\Sigma \rangle \langle J'\Omega'|\alpha_{Zx} + i\alpha_{Zy}|J\Omega \rangle + \langle \Lambda'S'\Sigma'|\mu_{z}|\Lambda S\Sigma \rangle \langle J'\Omega'|\alpha_{Zz}|J\Omega \rangle,$$
 (10)

where "has been dropped on the lower state quantities. Only one of the three terms on the r.h.s. of Eq. (10) is nonzero for any allowed transition matrix element. The electric dipole transition matrix is formed in the same basis as was the Hamiltonian matrix, that is, the Hund's case (a) basis. Selection rules in this basis are $\Delta S = 0$, $\Delta \Sigma = 0$, $\Delta \Lambda = 0$, ± 1 , and $\Delta J = 0$, ± 1 . Because $\Delta \Sigma = 0$ and $\Omega = \Lambda + \Sigma$, the selection rule $\Delta \Omega = \Delta \Lambda$ is obtained.

For ${}^2\Sigma^{-2}\Pi$ transitions the last term in Eq. (10) is always zero. Because we are interested in ${}^2\Sigma^{-1}\Pi$ transitions, we set all matrix elements of the type $\langle {}^2\Sigma|\mu_Z|^2\Sigma \rangle$ and $\langle {}^2\Pi|\mu_Z|^2\Pi \rangle$ equal to zero. This approximation is valid because these matrix elements are much smaller in magnitude than the $\langle {}^2\Sigma|\mu_Z|^2\Pi \rangle$ matrix elements and because they enter into the line strength only through weak mixing between the ${}^2\Pi$ and ${}^2\Sigma$ wavefunctions. The relative phases of the matrix elements are determined by following the prescription of Hougen⁴ and Whiting and Nicholls. Following the suggestions of Whiting and Nicholls, we have normalized the line strengths so that

$$\sum_{SJ'} S_{J'J'} \approx 2(2S+1)(2J''+1) = 4(2J''+1). \tag{11}$$

Table 4 reflects this choice of normalization.

The line strength $S_{P,P}$ is formed in intermediate coupling by taking $|\langle u|\mu_Z|l\rangle|^2$, where $|u\rangle$ represents the eigenvector of the $^2\Sigma$ state, $|l\rangle$ represents the eigenvector of the $^2\Pi$ state, and μ_Z now represents the 6×6 transition matrix. However, the eigenvectors formed during diagonalization of the energy matrices are in the Kronig transformed basis; therefore, the dipole matrix must also be transformed into this basis. Using the Kronig transformed wavefunctions given by Eq. (1), the Kronig transformed dipole matrix elements have the form

$$\langle J'\Omega'\delta'|\mu_{Z}|J\Omega\delta\rangle = \frac{1}{2}(\langle J'\Omega'|\mu_{Z}|J\Omega\rangle + \delta\langle J'\Omega'|\mu_{Z}|J-\Omega\rangle + \delta'\langle J'-\Omega'|\mu_{Z}|J\Omega\rangle + \delta\delta'\langle J'-\Omega'|\mu_{Z}|J-\Omega\rangle)$$
(12)

and, if we let μ_z^x be the matrix of the $\langle J'\Omega'\delta'|\mu_z|J\Omega\delta\rangle$,

$$S_{J'J'} = |\langle J'\Omega'\delta'|\mu_Z^K|J\Omega\delta\rangle|^2. \tag{13}$$

To our knowledge, the ${}^{2}\Sigma^{-2}\Pi$ transition matrices have not been published elsewhere. We show them for the P, Q, and R branches in Tables 5-7. The twelve branches correspond to the four 3×3 blocks in the Kronig basis as follows:

$$\begin{pmatrix}
a \leftarrow a & a \leftarrow s \\
PP_{11}, RR_{11}, PP_{12}, RR_{12} & Q_{11}, PQ_{12} \\
s \leftarrow a & s \leftarrow s \\
QQ_{22}, RQ_{21} & PP_{22}RR_{22}, QP_{21}SR_{21}
\end{pmatrix} (14)$$

For the Q-branch lines, the only nonzero 3×3 blocks are those for $a \leftarrow s$ and $s \leftarrow a$, while, for the R- and P-branch lines, the only nonzero 3×3 blocks are those for $a \leftarrow a$ and $s \leftarrow s$. When forming line strengths, the appropriate block must be substituted into Eq. (13) for μ_{Σ}^{E} . For

Table 5. P-branch 2x-30 transition matrix.

Hati	1) Mmd's case (a) basis	basis.					11) Kro	II) Kronig basis					
	2/4- N	$^{2\mathbf{E}_{-1}^{-1}}_{I_{2}} ^{2\mathbf{E}_{-1}}_{I_{2}} ^{2\mathbf{E}_{-1}}_{I_{2}} ^{2\mathbf{E}_{1}}_{I_{2}} ^{2\mathbf{E}_{1}}_{I_{2}} ^{2\mathbf{E}_{1}}_{I_{2}}$	2/1-12	2112z		2/83/2		28.3/E	2/182	2 1 2 2 1 2 2 1 2 2 1 2 2 1 2 2 1 2 2 1 2 2 1 2 2 1 2 2 1 2 2 2 1 2 2 2 1 2	21.12		28.5
2H-3/2	<u>°</u>	. •	7	э	0	/	21/2	<u>°</u>	•	7			
21,-12	•	•	9	Ħ	•	•	21/2	•	•	н		0	
7/1-1/	7		9	0	M I	9	252,	¥		•)	
2/13,	٥	ř	•	0	0	7	22,72				•	*	•
7/182	•	٥	Ħ	0	9	c	21,12		0		7	٠	0
ر. *ارا	•	•	•	7	•	0	283/2				7	٠	•
	•		-			\		•					•
•	[·] · · ·	(2/2)(1/2)(1/2)	1, _{1,1}										
	₹ F	1, 10(1, 10(1, 1)	(<u>1</u> 2) 1/2										
	Viere	where p Siprisp. - 1	1 - [^41 1]	_									
		•											

Table 6. O-branch 25.-211 transition matrix.

Head.	1) Hund's case (a) basis	basis					11) Kra	ii) Kronig basis	_				
	2/1-1/2	24 ₋₁ / ₂	2H-1/2 2H-1/2 2E-1/2 2E1/2	21,12	#1/2	2/1/3		2,12	201, 2E3/2 2E1/2 201/2 201/2 201/2	213/2	2,12 1,12	2/182	20,2
21,21/2	0	•	7	•	9	/	283°	\			<u>.</u>	•	9
1,-12	ė	•	•	7	•	0	281,2		0		>	•	•
21,1/2	Y	•	•	٠	•	•	21,12				•	*	•
7E1/2	0	-	•		•		21,12	•	•	•			
2/102	°	•	>	•	•	•	281/2	°	•	†		0	
283/2	ر ر	•	•	.	•		24.3/2	بر	•	•		1	
	. :	0(17 ₁ -1)		1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1		•		•					
		3	<u> </u>	•									
	v · I	270740 24041 7 Zely	v = \(\frac{(10^1/2)^2(2341)}{23(341)}\) \(\frac{\pi}{23(341)}\) \(\frac{\pi}{	· •• •									

Table 7. R-branch $^2\Sigma$ - $^2\Pi$ transition matrix.

	A) Mand's case (a) basis	Fasia					11) Pro	11) Kronig basis					
	212-3/2	29-1/2	74-312 24-12 2E-1/2 2E1/2 2A1/2	2/132	2/182	2,482		2 18 3 1/2		281/2 251/2 251/2 281/2	2 L1 .	381,2	283/2
2,12	,	•	,	G	¢		7E3.	′	•				
		•	•	•	•	•	~	•	•	>			
11.	•	•	•		•	•	2112	•	•	7		0	
21,-1/2	•	•	•	•	*	•	21,12	,	7	•)	
2F1/2	9	*	•	د	•	, x	21,12				•	•	•
2/1/2	•	•	H	•	•	•	281/2		C		•	•	•
2/5/2	<u>,</u>	9	•	>	•	•	283,8)			•	6
	A	1002 1001-11	y - [(1-1/10)(10+1/10)] · y	.4		•							
	z • į!	1007-17 1007-17	$z = \frac{(10^{1}/2)(10^{-3}/2)}{2(101)}^{1/2}$ where $v_{\Delta} = \sqrt{ v_{\Delta} ^2 + 1}v_{\alpha} = 1$.4 -									

example, to generate the Q_1 line strength we form

$$S_{s}^{2} = |\langle \psi_{s_{nin}}^{2} | (a \leftarrow s) | \psi_{s_{nin}}^{2} \rangle|^{2}.$$
 (15)

3. RESULTS AND DISCUSSION

Line strengths, Einstein A coefficients, intersities, and transition frequencies have been calculated for all branches of the $A^2\Sigma - X^2\Pi(0,0)$ OH spectrum through J=15.5 at 240°K for atmospheric applications and through J=40.5 at 4600°K for high temperature applications. These calculated values are shown in Tables 8 and 9, respectively. The total band intensities (by summation of the individual lines) are $2.7948 \times 10^4 \, \mathrm{cm}^{-1}/\mathrm{atm}$ -cm at 240°K and $8.6863 \times 10^2 \, \mathrm{cm}^{-1}/\mathrm{atm}$ -cm at 4600°K. Line intensities are plotted at these two temperatures in Fig. 1. Some caution must be exercised in using high J data from Table 9. The spectroscopic constants 2 used here were determined from data 2 which included transitions through J=25.5. Although these constants allow prediction of that data to within 0.1 cm $^{-1}$ maximum error and a standard deviation of $\sim 0.03 \, \mathrm{cm}^{-1}$ (the hyperfine structure is neglected here), such accuracy cannot be expected for all lines between J=25.5 and 40.5. Uncertainties in the calculation of energy levels at these high J cause proportionally smaller uncertainties in the energy eigenvectors (wavefunctions) and in quantities calculated using the eigenvectors (line strengths, Einstein A coefficients, and intensities).

Chidsey and Crosley' list p_*^{*} through J=35.5. We have extrapolated p_*^{*} for J=36.5. 40.5. Although the dependence of p_*^{*} on J is quite linear for P, Q, and R transition probabilities between J=25.5 and 35.5, extrapolated p_*^{*} used to calculate Einstein A coefficients and intensities at higher J must be used with caution.

Although we list four digits for Einstein A coefficients and intensities in Tables 8 and 9, the absolute uncertainties of these quantities cannot be less than 1%, because German's t^{10} t^{10} has a 1% quoted uncertainty and because Chidsey and Crosley's t^{10} are quoted to three significant digits. The relative uncertainties are limited by the relative accuracy of the t^{10} and the line strengths. These relative uncertainties should be less than 0.5% for t^{10} less than 25.5 and are probably less than 1% for t^{10} less than 35.5.

Our calculated line strengths have been checked for accuracy by comparison with values calculated using Earls¹⁶ algebraic formulas (these formulas are equivalent to those of Kovacs¹² for $^2\Sigma^{-2}\Pi$ transitions). When centrifugal and higher order distortion and $\Sigma^{-1}\Pi$ interactions are

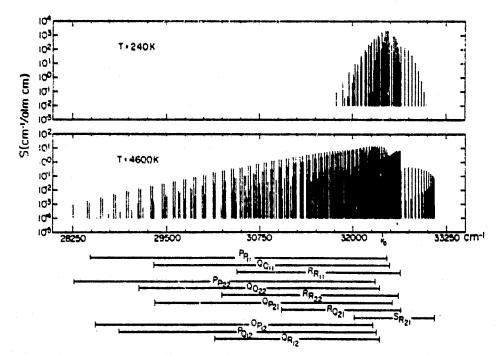


Fig. 1. Line intensities and positions for the $A^2\Sigma - X^2\Pi(0,0)$ band of OH. Only lines with intensity greater than 10^{-5} the intensity of the strongest line have been plotted.

Table 8. Line parameters for the $A^{(\Sigma-\lambda)^2}$ H(0,0) OH band at 200°K.

	Vacuum Cm-1	mer gy	Transition	Frequency	Mavelength	Intensity	Intensity ca/molecule	Einstein A	Line	
						T = 240	O K			
_		- 226-	211	30975.003	3227.4696	2.1106-11	16-3256-31	3:9116.03	2.203416-81	-
3.5 16.	2	7		31004.643	3216.00%	6.6496-10	2.1746-29	4.547K+03	2.310606-01	~ (
		,		260-16135	3583-8355	1.72%-11	2-X4-6-	9-36K-63	10-3707X	,
_			2111-5	11 396 960	11 61, 6942	4.514F-A6	2-136-26	7.106.01	2.724045-01	•
			2 2015.5	316.50.396	1177-ABS	1.6236-00	S. 36 W-26	7.000	1.594745.01	•
			1 2115.5	31463-010	1177.1155	7.4796-11	2,446-34	1.2815.06	7.573446-11	~
-	2453	5	1 2(10.5	31490.667	31.73.61.93	9.3736-05	3.064-24	9.00 × 003	2.903356-01	•
_	1371	365	P 2 2114	31540.460	3169.6197	4.5176-11	1.4716-27	3.0676.05	1.491966.01	•
Ξ	1367	7	9 1 2114.5	31543.628	1169-2961	2.3336-09	P. 62.E-29	1.476.04	7.915666-01	=
Ξ.	1361	•	P 1 1115.5	31573.139	3166.3329	5,7546-11	1.006-27	3.15%.05	1.62750€.01	=
	28 55		P 1 21 9.5	31595.731	3164.0607	1.1996-13	3.5956-23	1.0%	3, 103666-01	21
_	30.1		2 2(13.5	31617.484	3161-2998	1-0336-06	3.376-26	3.1766.05	1.309176.01	2
Ξ	717	. 978 .	0 1 2013.5	31623.425	3161.5977	9-3610-9	1.566-27	1.71E.F.	10-366206.0	=
-		.759	P 1 1111.5	-	3150.5892	1.3376-06	4. 371E-26	3.1726.05	1.526952	15
	1631		1 21 1.5	325.98.028	3154.6534	1.0406-62	3.42K-22	1.3526.1	3.32%2E-01	•
_	334	- 225	2 2012.5	31691.411	3154.5157	1.9266-15	6.29%-25	3.26Æ • 05	1.20593€+11	11
_	334		9 1 2(12.5	31694.228	3154.2353	1.2006-66	4.1076-26	2.1056.14	10-3/021/19	=
Ξ	3311	- 206-	1 1113.5	31726.769	3151.0001	2,5476-65	1.3366-25	3.20%.05	1.426156.11	=
_	2001	.525	2 2111.5	31762.412	3167.4639	2.9206-64	2-XX-5	3.3626.05	1-102176-01	2
=	2001	. 525 .	1 2011.5	31765.021	3147.2054	2,2356-15	7.3166-25	2.35% • 84	9.24697E-01	2
246	901	.726	1 21 7.5	31701.413	3145.5021	1.176-12	2.6416-21	1.6956.1	3.57791E-01	2
-	9712 S	į	-	31799.454	3143-7472	3.954 E-04	1.29X-23	3.40% . 15	1.32510€ • 11	22
_	9652 5	- 612.	P 2 2(10.5	31636,363	3146.7425	3.57 £E-03	1.1766-22	3.4496.85	1.177006.11	z
_	2,51	- 613.	1 2011.5	31032.700	3140.5061	3.1956-04	1.04.23	2.0016.00	1. 11 M. LE-11	X.
_	5 2413	- £19.	1 1111.5	31069.446	3136.8927	4.9A4E-13	1.6366-22	3.516€.85	1.224226.01	ž
	1170	:	P 1 20	31069.705	31 36. 0672	5.0036-01	1.6346-20	2.1626.04	1.144106-01	2
	202	151	2 21 9.5	31695.203	3136.3516	3.5256-82	1.15×-21	3.5326.05	4.727275.00	2
	502	151	21 9.5	31497.463	31.34-1373	3,7246-63	1.2106-22	3.35% · 04	1.045306.00	٤
-	2015	136		31936.701	31 30 . 27 47	5.17%-62	1.661E-21	3.6286.05	1.12265€ • 01	#
	124	=	P 1 20 5.5	31954.665	3120.5267	2,4746.08	1.132-20	2.01%.04	4.119306-01	ī
	1691	.23	P 2 21 1.5	31957.040	3126.2934	2.7796-01	12-3:11%	3.6116.05	L. 66.126E • 81	~
	1693		0 12 1 0	31959.000	3120-1015	3.5306-12	1.154-21	4.1776.14	1.1105PE+00	22
-	3	Ţ	9 2 2115.5	31963.173	3127.6939	3,016-09	9.0426-29	5.54% • 15	3.120306.01	3
	-		P 1 11 9.5	32001.569	3123.9476	4.1726-61	1.364-21	3.745€+15	1.021066.01	35
	1367	.617	1 12 2 .	32015.547	3122.5733	1.740€ • 60	8.60¥-20	3.640€ • 85	7.59%te.	*
-	4111	.419	0 2 2614.	32315.910	3122.5417	1.4366-00	2.75 Æ-27	5.71X .05	2.924916.01	2
	1367	3	1 21 7.5	32017.324	3122.4338	2,7126-61	8. 876E-21	5.120€ +04	1.200926.00	2
-	4377	;	A 1 2014.	32019, 325	3122.2887	1.4346-19	4.69E-29		10-321521-5	S
	9		P 1 21 4.5	32035.476	3128.5458	9.645€.11	3.154-19	3.7616.01	4. 12924E-01	;
•	1517		P 2 1(15.5	32042.018	3119.9173	1,000 6-69	6.1756-29	1.1956.04	5.574136-01	7
-	4351	2	1 1115.5	32845.433	3119.6649	1.1266-17	3.3546-27	5.5776.05	J. 12011E . 01	?
	1351		P 1 11 1.5	32063.69	3117-0005	2.7516.68	1.97.1-28	3.169€+15	9.19253E · BB	7
~ •	3986	- 651	9 2 2113.5	32864.936	3117.7762	1.9396-16	6.3416-26	5.164E+85	2.7210EE+ B1	;
-	3846	Š	R 1 2013.5	32868.126	3117.4571	3.7426-68	1.224-27	1.156.11	5.401506-01	Ş
2	1077	•	~	32879.768	3117.2084	6.575E+00	2.0045-19	3.7376.85	- 10	3
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1.2676.06 5.736.05 6.026.05	1.23E-64 3.70SE-05 3.997E-05 7.991E-04	5. 600 C. 15	7,357,00 6,226,05 1,711(-06, 3,036,05 1,377,05 2,666,06	6.13% 65 2.0716.05 3.046.05 1.114.65 1.114.05 6.34% 05	\$6.346.2 \$6.346.2 \$6.346.2 \$6.346.2 \$6.346.2 \$6.346.2 \$6.346.2	
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2.939E-06 2.396E-06 3.639E-05	3-2-9E-01 3-2-9E-01 1-1-4-9E-01 0-3-1-E-01	5.576-65 5.576-65 9.696-81 3.2926-81 1.096-85	6.336.61 7.1266.64 6.0466-04 2.0466-04 2.1276.62 1.0146.42 2.7646-62	8.906-63 2.406-93 2.6115-02 3.2916-62 8.7616-81 2.4076-02	3.2956-03 3.2956-03 3.2956-03 3.2956-03 1.756-61 1.756-61 3.2956-03 1.756-61 1.756-61 3.2956-03 5.2956-03	7.496F-61 1.654F-61 1.914F-60 5.512F-60 1.914F
3115.0709 3115.7663 3113.37.22	3112,1992 3112,1992 3112,6644 3112,6644	3110-2196 3109-3416 3109-3617 3107-5617 3106-5559 3106-566	3106.6542 3106.0141 3105.6743 3135.4225 3103.3503 3102.3503 3102.3503	11 02 11 45 11 02 11 12 11 02 11 12 11 02 11 12 11 02 11 12 11 02 11 12 12 02 11 12 12 02 11 12 12 02 11 12 12 02 12 12	100 - 100 -	595) 1686 6361 1686 9952 1686 6566 7686 6267 7686 7259 7686 7259 7686 7259 7686 7259 7686 7259 7686 7259 7686
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A Line	strength		6.529126.00	1.02996.00	10-301016-6	2.7525X·2	4.5555E+20	4. 37233E-01	1.0 307E.	1.17366.01	1.01636E+80	1.263526.01	3.106266.00	1.11315E.DO	1-052666-01	10-3599999	1. 22124E+ 00	0. 41497E • 00	2.356456.00	1.336516.00	6.314.06.00	1.41.79E+88	1.45015€+00	4.25111E+00	1.564676.88	2.25.621E+00	2.2727%•00	3. 2065 Me • 00	10.34(925.1	6.1617X·0	9. 80 80 2E-01	10-369929-9	1.32577E.01	5.17621E+#6	1.273055.00	1.546406.00	6.779306-01	1. 2250XE+01	6.104076.00	1.34301E+	2.52565E+00	1-12-106-01
Einstein A	9ec_1		5.04X . 05	7.3156.1	4.656E.P.	1.976.15	5.4666.05	1.1286.05	9.221E+h	6.3576.05	5.99E.E.	6.31Fe.05	5.17Fe . 05	7.966.6	6.214.15	1.2126-15	1.1026.05	6.02×105	0.611E+05	1.6166.05	5.78E·15	1.711E+05	2.646.05	5.150E • 05	5.735€.65	4.110E · 15	2.650€+05	2.31X·15	2.650€+85	2.490€+85	1. EX.15		2.71E·15	2.626.15	1.5576415	1.464.18	1.2716.14	2.756.15	2.717E-65	1.226.05	1.0366.05	2-7916-05
Intensity	cm/molecule	240 K	1.40×-17	2.19%-10	2.XX.S	2.54 X-17	2.2846-17	0.62X-10	5.1376-10	1.XX-1	2.00 K-19	3.4506-10	5.41 X-17	1.1916-11	1.114-17	1.17.	4. M.E-11	2.79K-17	6.73%-17	1.12%-17	5.20 X-17	1-38%-17	2.50 K-17	7.30K-17	1.476-17	21-3211.9	1.13%-17	6.066-10	1.37%-27	3. C. K-10	2.004-17	1.0X-17	3.11.X-26	1.01%-10	2.1946-17	2.6416-17	7. 6376-29	5. e7 iE-25	2.638-19	1.1346-17	2.11.K-17	12-306.0
	cm-2-atm	T = 2	4.20%**	6.7116.01	1.6416.00	7.776E+12	6.3016.12	2.6378.02	1.5716.62	2.559ۥ81	1.501E+11	1.055€+12	1,679€+83	3.642€+01	3.4126+62	2.7215.02	1.250€+02	1.5-16-12	2.060E+03	3.4526+82	1.6216.63	3.9826.02	7.6476.02	2.233€+03	1.369€+63	1.9616.03	3.4726+62	2.1986.12	4.2056-61	9.300E-01	1.575E+12	5.771E.02	9.616E-17	3,110€+01	6.715E+12	1.176€+12	2.396E-19	1.7956-15	1.156E+11	3.474.6.62	6.475€.62	2.723E-84
Vavelength	STP angatroms		30%.3741	3147.2735	3109.2001	3019.0717	3119,1553	3611.0156	3169.7770	3049.7342	3997.465	3007.34.22				_	363.26	3013.27%	201.67	3101.6259	3011.5479	3616.2376	3000.0117	3879.9568	3070.4762	3673.4520	3077.0346	3074.3765	3073.4310	3072.2072	3072.0606	3072.0131	3671.1483	3171.455	3171.356	3070.31.62	3071.256	3069.6700	3069-1046	3661.1053	3061.7061	3060.6120
Frequency	_		32349.158	32354.203	32354.112	32356.410	32354.581	32 354.997	32355.411	32355.050	32379.485	32300.917	32396.659	32402.122	32483.485	35115.452	32422.526	32423.579	32441.541	32440.940	32441.001	32455.599	32457.981	32450.568	32474.170	32474.523	32409.364	32517.473	32531.710	32540.433	32541.901	32542.448	32551.652	32551.679	32559.631	32566.452	32561.229	32567.242	32572.485	32576.504	32577.557	32570.555
Transition	•		19.5 12 2 1 1	0 4 1 26 3.51	4 F 2 11 7.51	1 1 2 26 1.51	0 0 2 2(2.5)	0 R 1 2(1.5)	0 R 1 2(2.5)	0 0 1 11 7.51	0 F 2 11 6.51	9 9 1 11 6.51	15.2 11 1 4 4	0 F 2 16 5.51	19.5 11 1 0 0	R 2 21 .51	0 + 2 1(4.5)	9 9 1 11 4.51	15.1 11 1.41	0 F 2 11 3.51	0 0 1 11 3.51	R R 2 26 1.51	1 - 2 11 2.51	4 4 1 16 2.51	0 - 2 11 1.51	1 4 1 16 1.51	R R 2 26 2.51	R 4 2 21 3.51	4 R 2 2(14.5)	R R 2 2(4.5)	R 0 2 16 1.53	2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	251	8 12 2 1	~	K 1 16 2	9 2 1115	~	9	0 2 11 3	R 4 1 11 3.51	A 4 2 2(11.5)
Louer energy	Vacuum cm-1		+29.275 -	- 523-529	1029-092 •	167-781	200.769 •	107.491 -	200.769 •	1029-892 •	769.216 -	769.216 -	• 674.80	544-119	544.119 •	126.449 -	155.900 -	355.988 -	- 000.0	202.370 •	202.378 +	167.751 •	13.920 -	- 926-69	• 950-	• 951-	- 110-602	+59-429	- 321.365 -	- 401-109		- 000-	30-1-676 •	124.525 +	• 614.50	• 674.80	4341.846 -	1344.522 -	1077.654 -	•	i	2000.525 +
<u>.</u>	-		3.5	1.5	7.5	5-1	 	1.5	2.9	7.5	5:3	5:9	5. %	5.5	5.5	ŵ	6.5	5:		.5	3.5	1.5	5.2	2.5	1.5	5.1	5. 2	3.5	16.5	6.5	5:	•	13.5	\$.5	5.2	5.2	15.5	12.5	6.5	3.5	3.5	11.5
-			3.5	÷:\$	6.5	1:5	5.5	5: 2	3.5	1.5	5.5	6.5	1.5	4.5	5.5	1.5	3.5	4.5	Š.	5.5	3.5	5.5	1.5	5.5	S.	1.5	3.5	;	15.5	•	5:	s,	_	6.5	S.2			•	7.5	7.5	4:5	12.5

222122	2333555555	:	
7.11504.0 1.021194.0 1.021194.0 1.020194.0 1.020194.0			
2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.		25 27 27 27 27 27 27 27 27 27 27 27 27 27	
1.626 6.226 5.337 1.337 1.346 2.596 1.296 3.296	1.3366.2 1.3366.2 1.3366.2 2.3366.2 2.346.2 2.346.2 3.477.2 3.		
366.2859 367.2339 367.736 367.736 367.736	1067- 966 1067- 7610 1066- 1109 1066- 1	366.1316 366.1316 3663.926 3663.926 3663.926	10 1 2 1 2 1 2 1 2 1 2 1 2 1 2 1 2 1 2 1
32562.819 32562.819 32562.653 32567.372 32560.563	3291.026 32991.109 32693.109 32663.023 32665.046 32615.746 32615.746 32617.446	32625-117 32625-137 32627-384 32627-384 32637-384 32631-345 32631-35 32631-35 32631-35	516.15.15.15.15.15.15.15.15.15.15.15.15.15.
1967-617 2650-279 2650-279 2650-279 2693-290 2895-393	245.125 241.526 241.526 241.53	2613.613 - 2015.014 - 2015.014 - 2015.016 -	
		v. – –	

249.0 K 2.79460E+04 CH-2 ATH-1 AT T = THE INTEGRATER INTERSITY FOR THE BAND IS

9-1393E-16 CF/HELCHE

8

Table 9. Line parameters for the $A^{1}\Sigma^{-}X^{2}H\theta_{0}$ 0) OH band at 460°K.

	 3
. Line atrength	3.363866.01
Einstein / sec ⁻¹	9.645E.D.
Intensity ca/molecule	2.7126-20
Intensity ca-2-ata-	4.3276-62
Wavelength STP angstrome	3370.0956
Frequency	23507.766
Transfition	 2632
Louer energy	1231
**************************************	5 32.

3	;	7	3	2	S	3	\$	\$	5	3	Ş	3	3	2	3	3	5	3	3	3	3	Z	Z	2	2	Z	Ľ	2	2	2	£	8	=	≃ :	2 ;	£ 1	6 2	3 2	3	E	*	F	2	7	\$	F	*	7	7 :	ļ	
5.241996-01	1.54136E-01	•	7.54726:01	2.964176-01	3.52 Pt. 26-01	7. 79674E.BI	3.2030200	5.3110 E-01	1-5246	3. 3226 Me. 01	7.347625.01	2-942346-01	3.5351X-01	7.537356.61	3.10.206.01	5.351196-01	1.655156-01	3.227.90	7.14791E . B.	3.000546-01	3.54611E-01	3.0044-6.01	7. 3370FC-01	5. 39925E-81	1.698136-01	3.124300.01	2.904695.01	6. Well X:01	5.1567ef-01	3.027976-01	1.73006-01	1-3613X-11	7.136366-31	3.02505	3-91-96-31			1.77666-01	6-74-27E+21	2.925746.01	3.049046-01	3.501696-01	6.9306E.01	2.70470E·01	5.617HE-01	1.92666-01	2.02630E.BI	3.0160%.00	6.5403X6.01	7-2/5/100-7	
1.3456.03	9.1416.02	1.796.04	1.141E+1	2.4676.12	3.73X . 02	6.2716.05	1.07% .05	1.56K·13	1.0366.03	9.06E.P	1.976.04	3.2775.02	4.76% . 02	0.1006.00	1.10% .05	1.0176.03	1.1705.13	1.106.15	1.1976.05	4.1656.02	5.9346.82	1-306-05	1.016.05	2.03.63	1.375.63	1-2166.05	1.4275.65	1.3956.05	2.3976.03	5.69E-82	1.5176.03	7.10K:12	1.2166.05	1-336C+05	\$. X25-1		2012616	1-7216081	1.5976.05	1.45%.05	6.1626.02	6.521E+02	1.116.65	1.6616.85	3.1116.63	1.9666.03	1.500E+05	%	1.0186.05	100 MM - 1	
3.09E-22	1-1926-21	2. % 1K-28	5.XX-21	1.63%-23	2-1916-23	1.7X	3.58-20	5.434E-22	1.67.21	3.74×-20	0.07 K-21	2.93K-23	4.2716-23	7.6966-21	2. K-3.5	1. W.K-22	2.3716-21	5.41%-20	1.416-20	5.07 K-23	7.05%-23	2-XX-2	1.23%-26	1.32 1.	3.2965-21	7.717-20	1.114-19	2.1926-20	1.93K-21	1.22 M-23	1.546-21	1-134-22	1.97 X - 28	1.006-19	22-245-5	1.33E-13	2-X4C-C	£.20€-21	3.31 16-28	1.5066-19	1.31%-22	1.77.K-22	3.03%-26	2.00K-19	4.03E-21	1.49%-21	2.050E-19	1.1716-21	22-30C-3	2-76 W-22	
6.2216-66	1.9026-03	4.0056-02		2.605E-15	3.9746-05	6.050E-03	6.2076-02	9.4676-54	2.6916-63	5.9616-12	1.4168-62	6.775E-65	6.0146-05	1.1966-62	9.0026-02	1.4216-63	3.703E-03	1.6456-02	2.2636-82	1. MIE-15	1-1266-04		1.9776-02		5.2506-13	1.2316-61	.TITE-E	3.4976-62	3-3060-63	1.3136-64	7.2526-63	1-1-1	3.1666-02	1.7336-01		10-3766	11-3036-4	0.090F-01	5.2066-02	2.4026-61	2-1996-14	2.03E-M	4.0306-02	3.3216-61	6.430E-13	1.3415-62	3.2016-01	1.0705-63	7.9656-12	\$ -302E-64	
3170.1514	3375.4736	3375.0017	3177.8716	3376.065	1169.2939	3360.6072	3362.7004	3362.6699	3364 - 56 36	3359.7905	3352.2211	3351.5459	3156.7196	3350.0644	3347.2091	3366.6661	3365.4732	3344.3984	3334.6684	1333.9970	3333-1716	3332.50%	3332.5000	3131.9017	3330.000	3329.61%	3310.3330	3316.0312	3317.7426	3317.2730	3316.6592	3216.5336	3315.0026	3315-6502	7177-9877	3300-1646	1234-1650	3382,0425	3702.2000	3301.0506	_	3300.7787	3300-1322	3241.6616	3291-1113	3	5	3207.4903	7156-7626	3285-0541	
29593.446	29612.505	2611.306	29650.270	29664, 322	236.17.85	29677.338	29729.391	29735.030	29748.962	29755.148	29022.396	29626.434	29435.587	29847.595	5117.99862	29071.096	29882.546	122.26062	29479.439	29915.414	54992.829	24190.793	2999.795	32.,50	30013.664	31124.126	36126.944	X129.684	30122.310	11115.611	70 142.147	162.291	30149.207	39153.139	100.25205	38638.433	11256.211	30260.236	30273.392	38277.329	39279.251	30207.234	38293.3.4	36378.941	39376.124	30391.910	•	38469.587	13616.665	36424.958	
₹	2126		2 2137.	-	~	-		-	**		2 2136	2136	2 1137	1 1637.	2 2130.	12 1	÷	=		-	2 1136	~	1136.	1 2029	-	3	N	N		2(34.		2 1135	1135	1129			717	2121-	2 2(3)	1 1120.	1 2133	- - -	1134.	2 2126.	~	1 2120.	-	26.5	2512 2		
291.624 - P	ğ. ; ; ;	571.723 - 0	•	136.737 -	3721.729 - 0	121.729 -	.392 •	· 161 ± .392 · ·	Se.173 -	4 . 192, [6]			2671.416 • •	.71.416	16692.271 - P	•	•	•	2155.000 - 00.22E		•	•		1.237 +	1.267 -	•	. 263 -		7	•	. 205.6	A.586 .	•	,			1.328 +	ģ	•	1 502.01	- 126	97.982 -	97.302 - 0	P. 371 -	90.378 -					06 02.331 0	•
35.5	.5 26.5	.5 33.5	7 % 17	37.5	30.5	30.5	~ ~	31.5	25.5	32.5	5 36.5	36.5	37.5	37.5	39.5	30.5	24.5	31.5	35.5	35.5	36.9	29.5	36.5	29.5	23.5	5.5	28.5	34.5	20.5	34.5	5.25	35.5	15.5	29.5	44.9	, , , , , , , , , , , , , , , , , , ,	27.5	21.5	33.5	20.5	13.5	34.5	34.5	26.5	26.5	5.02	21 5.75	35.5	11 5.55	32.5 32.5 106	

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A Line strength		5.93910E+01
Einstein sec-1		2.51Æ·05
I Intensity cn/molecule		1.996-19
Intensity cn-2-atm- T = 40		7-3501.6
Vavelengih STP angstrons	255.455 255	1232-25126
Frequency vacuum cm-1	######################################	0<0.226.96
Transition		1167
Louer energy vacuum cm ⁻ 1	244, 284 1848, 384 1848, 384 1	
*C *C		Ç

2222	2223	232323	=	1111111	1222223	**********
10-301415-5	3, 32,200C-01 5, 34,30C-01 3, 01,70C-01 3, 4,200C-01	2,7100C 1,100C 1			2. 2. 2. 2. 2. 2. 2. 2. 2. 2. 2. 2. 2. 2	
2. 20 E . 95 3. 91 JE . 95 2. 96 JE . 95 2. 96 JE . 95	1.57.60 6.61% 02 6.20% 03 6.20% 03	2.75K.05 2.35K.05 0.16K.07 5.05K.65 4.54K.03	20 3/10 6 20 3/10 6		2.59C 2.74C 6.74C 9.74C	
2. 3.74	22.22 2.22 2.23 2.23 2.23 2.23 2.33 2.3		25222522 25222533 25222533 252233	2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.2.	5.5%-19 5.5%-19 6.7%-26 6.7%-26 6.9%-19 7.00%-19	
5.366-82		7.13E-C1 7.13E-C2 1.36E-C2 1.36E-C2 1.36E-C3 1.3	2-36-6 2-24-6 2-24-6 2-24-6 2-24-6 2-32-6 2-32-6 2-32-6 3-3-6 3-3-6 3-3-6 3-3-6 3-3-6	2.136-02 1.796-02 2.1966-01 0.1746-01 5.276-01 6.5276-03	96-3211-1 10-3211-1 10-3211-1 11-3620-1 20-3250-3 96-3250-3 19-3152-6 19-3152-6	7.514-6 3.284-6 3.284-6 3.384-6 3.384-6 1.184-
3230.7272 3227.46% 3223.76%	3223-1510 3222-9190 3222-6733 3221-6421	321.5161 320.4151 3219.6562 3219.6562 3219.6363	. 26 9. 50 26 20 5. 50 26 20 5	3565-074 3265-095 3265 3265 3265 3265 3265 3265 3265 326	2265 2261 1587 3611 1587 3611 1587 3611 1297 3611 1297 3611 1297 3611 1297 3611 1297 3611	254, 2511 254, 254, 2511 254, 254, 2511 254, 254, 254, 254 254, 254, 254 254, 254
36 94 3.051 36 975.063 31 011.210	31016.576 31016.069 31026.957 31030.725	1107.324 1105.250 1105.250 1105.250 1107.215	11115.34 1115.34 11115.34 11111.47 11113.67 1113.67 1114.67 1114.67	11191.95 11191.092 11191.092 11191.092 11191.092 11191.092 11191.092	935-962 H 937-962 H 937-962 H 937-962 H 937-962 H 937-962 H 937-962 H 937-962 H 937-962 H	11 11 12 12 13 13 13 13 13 13 13 13 13 13 13 13 13
4474	***					7
- 200-2020	2823	*****	743.075 • 13016.946 • 13016.946 • 13016.797 • 13016.797 • 1456.207 • 1956.207	1966.207 1946.189 6011.926 6101.926 2154.973 2150.244	2136.244 2136.244 2136.244 2146.444 214	
21.5 22.5 14.5 15.5 27.5 27.5	-	*****				\$ 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5

Table 9 (Contd).

laber		::::::::::::::::::::::::::::::::::::::
Line strength		6.11.76.01 6.54976-01 2.90676-01 1.106376-01 1.106376-01 3.577916-01 3.73628-01
Einstein A sec-1		2. 3.65 2. 3.65 3. 3.65 3. 3.65 3. 3.65 3. 3.65 3. 3.65 3. 3.65 3. 3.65 3. 3.65
Y Intensity Cm/molecule 4600 K	0.100C-24 2.470C-16 1.167C	2.39ff-10 1.616-21 2.676-23 3.616-21 2.976-19 2.976-19
Intensity ca-2-ata-1 T = 46	1.29E-05 1.29E-06 1.2	3.015F.60 2.554E-03 4.274E-65 4.690E-01 3.151E-61 3.151E-61
Wavelength STP angstroms		3144.04.10 3142.2806 3147.724 3147.2854 3147.2854 3145.5814
Prequency vacuum cm ⁻¹ S		31736-423 31754-173 31759-733 31759-733 31762-412 31762-613 31761-956
Transition		
Lower energy vacuum cm ⁻¹	272.272 272	7474,376 - 14764,434 - 2466,434 - 2466,434 - 2466,525 + 2466,525 + 1368,726 - 6611,055 +

																																												•							
25	Z	Ž	Ē	252	2	ž	S.	ž	ž	25.	ž	3	Z	292	2	722	565	3	267	3		2	271	272	273	27	275	275	212	278	2	2	Z	~			Ž	Ž	23	203	X.	Ē	ž	S	£	£	ž	~	Į	£	Ĭ
4.20701E-01	2106.0	74.397E-0	3.930126-01	2.725326.01	1.077866.01	20124E-01	3.53391E • 01	1.667916-11	1.4635%-01	2.11050f.11	4.916106-01	1.224025.1	1.041106-01	3.724116.01	9.727276·80	1.015306.00	1.79666-01	3.3312%	1.654766-01	2.625066.01	5.110636-01	3.525796.01	4.735546-81	11.12.56.01	2.710726.11	4.119306-01	0.6682FE • 86	1.11050E+#0	J-12030E+01	6.479176-01	5.320526-31	3. 323146 • 01	10.356925.2	1.021065.01		6. 90 90 50 50 50 50 50 50 50 50 50 50 50 50 50	7.20007	5.125126-01	2.6000 % • 11	1.707546-01	1.32424-01	5.5741X-11	3.120116.01	9.192526.00	2.72106€+61	5.401506-01	2.42675€+01	6.521036.00	34162	24.7E-	3.151.9E-01
5.2176.03	3.4026.05	6.3766.13	1.11K·65	316	J. 1. 18 . 15	22	¥ :	2.41K.03	5.971E.03	1.20%-17	7.256 .13	3.5166.05	2.16Æ+1.	5.0106.05	3.532E • 05	J. 35 K . 11.	3.886E+82	5.37% . 85	6.1516.13	1.5166.05	1.20% • 13	5.216E · 85	2.75%·13	3.620€.05	1.4666.05	2.41%-64	3.6116.05	4.17E.X	5.54E.15	7.47/6.03	9.5866.83	5.405E+05	1.6272.005	3.7456.95		1.116.61	5.026.04	9.1056.03	1.51%.15	3.6146.82	3.7676.84	1.195€ •16	5.577E . 05	3.060€.05	•3090	366.	7.75	•	ZZE	•	1.2676.04
3.1645-26	4.27£5-10	3.6916-20	2.93Æ-19	1-36K-19	4. 00E-10	3.56-19	3.4606-10	2.336-21	12-3102-5	1.204-19	4.87E-28	4.6176-10	2. 351E-19	3.55X-18	4.24E-18	4.47 X-19	4.4698-23	4.10%-18	5.53%-28	1.0416-19	6. 3686-28	4.2305-18	3.3745-21	4. BBC-16	1.776-19	2.7516-19	4.29K-10	5.45E-19	4.7026-10	7.211E-28	1.25.20	4.9706-18	61-3604-2	5.0645-10	11-2/22-1	A. A155-21	6. 59EF-19	9.32%-20	2.426E-19	7.12X-23	3.1316-19	1.1506-13	5.74E-10	=	6, 1845-18	1.19X-19	3.3116-19	4.01%-10	•	12-361-9	1.346-19
5.1496-12	.A22E+1	5.P00E-62	4.610E-10	2.1436-01	216	5.416E-61	•	7	•			Ť	3.7526-11	5.569€.68	=		7.124E-05	= :	1.030E-62	2.938E-01	1.0156-01	6.7626.66	5.3046-63	7.7966.10	.136E-1	11-308-11	.155E+D	11-3/1/-01	7.6296.10	1.150E-61	1.3156-01	7.943E.E	3. 766E-61	4 .00 SE + 60		7-5625-01	1.851666	1.4076-01	3.46 DE-61	1.1366-64	4.996E-C1	11-3604-11	9.1716.00	8.194E.CB	9.0666.00	13-3136-1	5.2036-61	6.405€+88	55E • 1	•	2.1476-01
3145.1192	3	3143.4154	31 6 3 . 09 6 0	3141-1051		3146.5861	3139.1767	31 39.02%	3130.7061	3136.4137	3136.9747	3136.9927	31 36 . 06 72	31 36.5946	3134.3516	3134-1173	36.7.7.16	3133.5356	31 35 - 96 39	31 32 . 52 %	31.31.933.	3130.56.21	31 30 . 12 75	31 30 - 27 07	3129.4015	3120.5267	3120-2934	3120.1015	3127-6939	3127.3414	31 25. 27 03	3124.9267	16.59	3123.347		77.	3122.4838		21.454	=	z	Ξ	Ξ	2	2	2	2	7.28	17.053	15.11	3118-6769
31766.009	£	31063-318	3	31026.709	31030.303	31032.708	31046.260	31047.024	31053.222	31053.291	31 069 .613		31069.765	31472.576	31195.203	31097.463	31190.447	219.99616	27-01616	31913.035	31939.106	32933.869	31,935.264	31936.761	31943.638	31454.665	31 957 - 149	31 959.100	31963-173	31966.776	31 917 . 177	21991.469	27.06.15	105° 19826	766.61876	17616.677	32617.324	32019.325	32021.935	32624.345	32035.976		32045.433	32163.698	32864.966	32169.126	32070.095	32676.768	_		32092.604
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Intensi cm ⁻² -atm T =	3.2076+60	3.1%E+00 2.933E+00	1.2196-61	3.669€.10	3.212E+60	3.9366.68	1.501E-11	3,7786.60	6.417E-64	1.4625.00	16-356-91	1.7516.60	1.1976-14	5.5215.10	1.7666.00	2.124E+10	2.0316-61	1.346E+CE	5.261E+82	6-0196-60	10.3692.1	3.334C000	6.310F-40	2.505E-01	5.174E+68	6.932E.el	5.484E+11	6.31SE+10	1.660€ • 11	1.445ۥ10	1.262E-11	5.491E.88	1.475E+10	5.370E+60	10-300E-01
Wavelength STP angstroms		3877-6346				3874.3765			_		3872.0131	3671.1463	-							_		1961 4191	_		•	_	_	_		-		-		3865.9766	3065.3710
Prequency Vacuum cm ⁻¹		32449.384				32517.473	32535.634		32548.433		32542.484	12551.652	32552.137	32551.679							32576.584						•							32666.557	32612.991
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ignored, our line strengths are identical to those calculated using Earls' formulas. Ignoring these effects does not significantly alter the line strengths for main branch transitions but does lead to large errors at high J for the weaker satellite branch line strengths. Earls' forumlas (which ignore centrifugal and higher distortion and Σ - Π interactions) predict consistently smaller line strengths than those calculated by us (which include centrifugal and higher distortion and Σ - Π interactions). In the ${}^{Q}R_{12}$ branch, Earls' line strengths range from 93% of our line strength at J'' = 20.5 (~ 59% of our line strength at J'' = 40.5; in the ${}^{S}R_{21}$ branch they range from 78% of our line strength at J'' = 20.5 to 46% of our line strength at J'' = 39.5. Earls' formulas show intermediate error in this range for the other satellite branches.

Bennett's¹³ line-strength formulas, which include P^4 centrifugal distortion, may be expected to give much more accurate results. Based on our check of Earls' formulas, the line strengths for the ${}^5R_{21}$ branch should have the largest deviation. Bennett's formula predicts line strengths for this branch which are 1% higher than ours at J=1.5 and decrease to 6% lower at J=25.5; they increase to 2% higher at J=39.5. Examination of the eigenvectors shows that the P^6 term and, to a lesser extent, the ${}^2\Sigma^{-2}\Pi$ mixing can contribute an effect of a few percent to the satellite bands line strengths. Thus, the dominant effect in the deviations from Earls' formulas is due to the centrifugal disortion, which is relatively large in a light molecule such as OH.

The conclusion to be drawn from these comparisons is that both the algebraic formulas or our method give accurate main branch line strengths at all experimentally observed J values. Earls' formulas lead to significant errors at high J in the satellite bands. Bennett's formulas give acceptable satellite branch line strengths for most work. However, when the highest accuracy is required, the line strengths from Table 9 should be used.

When our Einstein A coefficients are normalized to the same relative value as those of Chidsey and Crosley,³ the two sets of values differ by at most ± 2 in the last decimal place. As with line strengths, these differences become inportant only at high J in the satellite branches where many of the relative Einstein A coefficients are quoted to only one significant digit by Chidsey and Crosley. An additional advantage of our Einstein A coefficients in Tables 8 and 9 for quantitative spectroscopy is that they are absolute rather than relative values.

Although we have chosen to present OH line parameters for temperatures of 240 and 4600° K, our computer program can generate $A^{2}\Sigma-X^{2}\Pi(0,0)$ band line parameters for any temperature. Table 10 may be used with Tables 8 or 9 and Eqs. (3) and (5) to convert line intensities from these temperatures to any temperature in the 200-6000°K range. Intensities so determined should have the same accuracy as those in Tables 8 and 9. Band intensities at temperatures other than 240 and 4600°K may be calculated by summing the individual line intensities at the desired temperature. Simpler approximate procedures which directly convert from a band intensity at one temperature to a band intensity at another temperature such as Eq. (7-126) in Penner, 6 give errors of approx. 15% when band intensities at 240 and 4600°K are compared.

Table 10. Rotational partition function On vs temperature T.

T (K)	Q _N	T (K)	Q.
200	26.71	2000	298.23
240	32.24	2500	375.27
296	40.17	3000	454,19
300	40.75	3500	533.18
500	70.00	4000	612.43
750	107.35	4300	693.14
1000	145.09	4600	709.26
1250	183.08	5000	773.98
1500	221.28	5500	855,28
1750	259.67	6000	936.88

Acknowledgements—This research was supported in part by NASA/Langley contract NSG 1465. We thank J. L. Destombes and C. Marliere-Destuynck for kindly providing mole-blar constants which predict the line positions better than those previously published, and I. L. Chideey and D. R. Crosley for kindly providing tables of P\$% prior to publication. Acknowledgement is made to the National Contex for Armospheric Research, which is spensored by the National Science Foundation, for computer time used in this research.

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ERRATUM

The following corrections should be made to Goldman and Gillis¹.

In Eq. (3) and (7) the expression (2J'+1) should be 4(2J'+1).

Consequently, all intensities in cm⁻¹/atm-cm and cm/molecule in Ref. 1 should be multiplied by 4. All other line parameters remain unchanged.

1. A. Goldman and J. R. Gillis, JQSRT 25, 111 (1981).